

THE EIGENSYSTEM OF THE EULER EQUATIONS

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In this document I list the eigensystem of the Euler equations. The formulas are taken from [1], Chapter 3, section 3.1. The Euler equations can be written in conservative form as

$$\frac{\partial}{\partial t} \begin{bmatrix} \rho \\ \rho u \\ \rho v \\ \rho w \\ E \end{bmatrix} + \frac{\partial}{\partial x} \begin{bmatrix} \rho u \\ \rho u^2 + p \\ \rho uv \\ \rho uw \\ (E + p)u \end{bmatrix} = 0 \quad (1)$$

where

$$E = \rho\varepsilon + \frac{1}{2}\rho q^2 \quad (2)$$

is the total energy and ε is the internal energy of the fluid and $q^2 = u^2 + v^2 + w^2$. The pressure is given by an equation of state (EOS) $p = p(\varepsilon, \rho)$. For an ideal gas the EOS is $p = (\gamma - 1)\rho\varepsilon$.

The eigenvalues are $\{u - c, u, u, u, u + c\}$. The right eigenvectors of the flux Jacobian are given by

$$R = \begin{bmatrix} 1 & 0 & 0 & 1 & 1 \\ u - c & 0 & 0 & u & u + c \\ v & 1 & 0 & v & v \\ w & 0 & 1 & w & w \\ h - uc & v & w & h - c^2/b & h + uc \end{bmatrix} \quad (3)$$

here

$$h = (E + p)/\rho \quad (4)$$

$$c = \sqrt{\frac{\partial p}{\partial \rho} + \frac{p}{\rho^2} \frac{\partial p}{\partial \varepsilon}} \quad (5)$$

is the enthalpy and the sound speed respectively. Also,

$$b = \frac{1}{\rho} \frac{\partial p}{\partial \varepsilon}. \quad (6)$$

Note that for ideal gas EOS we have

$$h = \frac{c^2}{\gamma - 1} + \frac{1}{2}q^2 \quad (7)$$

$$c = \sqrt{\frac{\gamma p}{\rho}} \quad (8)$$

and $b = \gamma - 1$. Hence, in this case the term $h - c^2/b$ in Eq. (9) is just $q^2/2$. The left eigenvectors are

$$L = \frac{b}{2c^2} \begin{bmatrix} \theta + uc/b & -u - c/b & -v & -w & 1 \\ -2vc^2/b & 0 & 2c^2/b & 0 & 0 \\ -2wc^2/b & 0 & 0 & 2c^2/b & 0 \\ 2h - 2q^2 & 2u & 2v & 2w & -2 \\ \theta - uc/b & -u + c/b & -v & -w & 1 \end{bmatrix} \quad (9)$$

where

$$\theta = q^2 - \frac{E}{\rho} + \rho \frac{\partial p / \partial \rho}{\partial p / \partial \varepsilon} \quad (10)$$

which, for an ideal gas EOS reduces to $q^2/2$.

Now consider the problem of splitting a jump vector $\Delta \equiv [\delta_0, \delta_1, \delta_2, \delta_3, \delta_4]^T$ into coefficients needed in computing the Riemann problem. The coefficients are given by $L\Delta$. For an ideal gas law EOS, after some algebra we can show that an efficient way to compute these are

$$\alpha_3 = \frac{\gamma - 1}{c^2} [(h - q^2)\delta_0 + u\delta_1 + v\delta_2 + w\delta_3 - \delta_4] \quad (11)$$

$$\alpha_1 = -v\delta_0 + \delta_2 \quad (12)$$

$$\alpha_2 = -w\delta_0 + \delta_3 \quad (13)$$

$$\alpha_4 = \frac{1}{2c} [\delta_1 + (c - u)\delta_0 - c\alpha_3] \quad (14)$$

$$\alpha_0 = \delta_0 - \alpha_3 - \alpha_4. \quad (15)$$

REFERENCES

- [1] Andrei G. Kulikovskii, Nikolai V. Pogorelov, and Andrei Yu. Semenov. *Mathematical Aspects of Numerical Solutions of Hyperbolic Systems*. Chapman and Hall/CRC, 2001.